

INVERSE MATRIX PROJECTIVE SYNCHRONIZATION OF NOVEL HYPERCHAOTIC SYSTEM WITH HYPERBOLIC SINE FUNCTION NON-LINEARITY

Smail Kaouache^{1,2} and Mohammed-Salah Abdelouahab²

¹Department of Mathematics
University of Mentouri Brothers, Constantine. Algeria

²Laboratory of Mathematics and their Interactions
Abdelhafid Boussouf University Center, Mila. Algeria

Abstract. In this paper, we investigate the inverse matrix projective synchronization (IMPS) of novel hyperchaotic system with hyperbolic sine function non-linearity. Recall that the studied system is generated from the modified Lü system. First, hyperchaotic attractors, symmetry, dissipation, equilibrium points and Lyapunov spectrum are the tools used to analyse this system. Moreover, this paper presents an active controller to achieve the IMPS analysis of the system. The main results are established by using Lyapunov stability theory, and finally numerical example and computer simulations are shown to illustrate all the main results.

Keywords. Hyperbolic sine function, Hyperchaotic Lü system, Active control, Lyapunov exponent, inverse matrix projective synchronization.

AMS subject classification: 34A34, 37B25, 37B55, 93C55, 37C25.

1 Introduction

Grace to the natural properties of chaotic and hyperchaotic systems, such as: sensitivities of initial conditions, boundedness and infinite recurrence, hyperchaotic systems have become good condidat for important applications in several areas such as: cryptosystems, secure communications, network signal transmission, electrical circuits and encryption [8, 12, 13, 23, 26], etc.

Chaos synchronization plays an important role in nonlinear science and must consider several aspects, such as physical systems [11], ecological systems [3] and biological systems [10], etc. For this, many different methods have been used to study the synchronization and stability of general uncertain systems, such as impulsive control, adaptive control, Adaptive fuzzy control, active control, prediction-based feedback control, sliding mode control [2, 5, 6, 17, 21, 22, 24] and so on. Different types of synchronization

have been proposed, such as complete synchronization, projective synchronization, co-existence of some chaos synchronization types, generalized synchronization, inverse generalized synchronization and hybrid synchronization [1, 4, 14, 16, 18, 19, 20].

On the other hand, most of the existing papers discuss the synchronization between two chaotic or hyperchaotic systems with quadratic or cubic non-linearity. Recently, there has been increasing attention to the synchronization of hyperchaotic systems which contain a non-linear term in the form of usual functions with different characteristics and behaviour. For example, in [9] an anti synchronisation of novel hyperchaotic system with its exponential non-linearity function has been investigated.

Also, the concepts of hyperbolic functions are analogues of the usual functions. For example, hyperbolic functions occur in the solutions of many linear differential equations such as the Laplace equations, which have attracted a great deal of interest in many areas of physics, including fluid dynamics, and special relativity. Owing to the previous discussion, a novel hyperchaotic system with hyperbolic sine function non-linearity generated from the modified Lü system [7] is investigated. In addition, the IMPS analysis of the system is achieved using active control method and Lyapunov stability theory.

The remainder of this paper is structured as follows: The next Section states the novel hyperchaotic system and its dynamic analysis. The Section 3 derives an active controller for IMPS of the identical systems. Finally, the Section 4 contains the conclusion.

2 A novel hyperchaotic system and its hyperchaotic attractors

Our novel hyperchaotic system generated from the modified Lü system [7] is given as

$$\begin{cases} \dot{x}_1 = b_1(x_2 - x_1) + x_4, \\ \dot{x}_2 = -x_1x_3 + b_2x_2, \\ \dot{x}_3 = \sinh(x_1x_2) - b_3x_3, \\ \dot{x}_4 = b_4x_1, \end{cases} \quad (1)$$

where x_1, x_2, x_3, x_4 are the state variables, b_1, b_2, b_3 and b_4 are positive real. In addition, we shall show that the system (1) is hyperchaotic when the parameters b_1, b_2, b_3 and b_4 take the values

$$(b_1, b_2, b_3, b_4) = (25, 20, 1.5, 7). \quad (2)$$

We recall that we can also find other values of the previous parameters to justify the existence of hyperchaos in the proposed system by using the bifurcation diagram and also by the Wolf algorithm [25].

All numerical simulations are performed by using the fourth order Runge-Kutta algorithm. For this numerical simulations, we take the initial values

of the novel hyperchaotic (1) as

$$\begin{aligned} x_1(0) = -0.0546, \quad x_2(0) = -0.0049, \quad x_3(0) = 0.0427 \\ \text{and } x_4(0) = 0.9094. \end{aligned} \tag{3}$$

The attractors of the system (1) are represented in Figure 1.

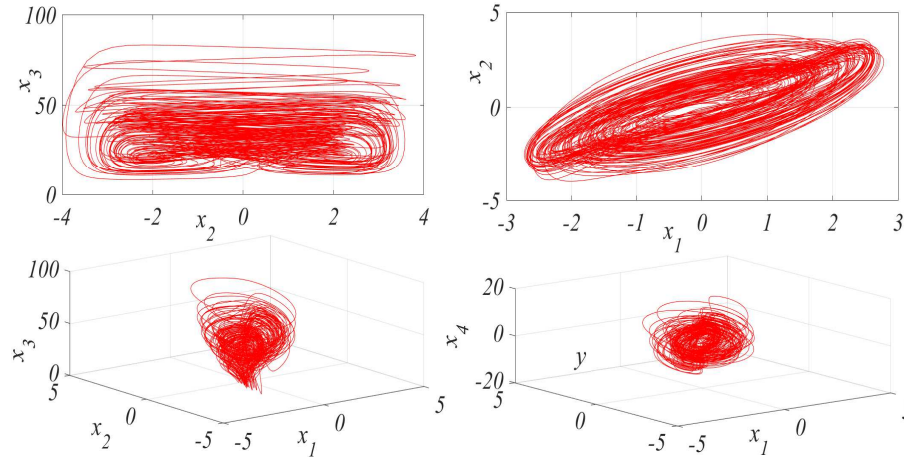


Figure 1: The attractors of the hyperchaotic system (1).

The proposed system has the following characteristics:

2.1 Symmetry

The system (1) remains invariant under the function:
 $(x_1, x_2, x_3, x_4) \rightarrow (-x_1, -x_2, x_3, -x_4)$. In addition, this system is symmetric with respect to the x_3 - axis.

2.2 Dissipation

The divergence of the novel 4 - D system (1) is

$$\nabla f = -b_1 + b_2 - b_3. \tag{4}$$

By using Liouville's theorem, we obtain

$$V(t) = V_0 \exp((-b_1 + b_2 - b_3)t) \tag{5}$$

For the system (1) to be dissipative, it is enough that $\nabla f < 0$. in this case, all the trajectories of the system tend to an attractor when $t \rightarrow +\infty$.

2.3 Equilibrium points and stability

Suppose that the system (1) has $P(x_1, x_2, x_3, x_4)$ as an equilibrium point, thus

$$\begin{cases} b_1(x_2 - x_1) + x_4 = 0, \\ -x_1x_3 + b_2x_2 = 0, \\ \sinh(x_1x_2) - b_3x_3 = 0, \\ b_4x_1 = 0, \end{cases} \quad (6)$$

As a result, the system (1) has only one equilibrium point $P(0, 0, 0, 0)$. The Jacobian matrix in P is

$$\begin{pmatrix} -b_1 & b_1 & 0 & 1 \\ 0 & b_2 & 0 & 0 \\ 0 & 0 & -b_3 & 0 \\ b_4 & 0 & 0 & 0 \end{pmatrix}, \quad (7)$$

and its characteristic polynomial is

$$P(\lambda) = (\lambda - b_2)(\lambda + b_3)(\lambda^2 + \lambda b_1 - b_4). \quad (8)$$

The eigenvalues corresponding to P are

$$\begin{aligned} \lambda_1 = -b_3, \quad \lambda_2 = -\frac{1}{2}b_1 - \frac{1}{2}\sqrt{4b_4 + b_1^2}, \quad \lambda_3 = b_2, \\ \text{and } \lambda_4 = -\frac{1}{2}b_1 + \frac{1}{2}\sqrt{4b_4 + b_1^2}. \end{aligned} \quad (9)$$

Here, λ_1 and λ_2 are negative real numbers, λ_3 and λ_4 are the positive real numbers. Then the equilibrium P is an unstable saddle point.

2.4 Lyapunov exponents

Here, we assume that the parameters b_1, b_2, b_4 remain fixed and just b_3 is varied in $[0.5, 1.8]$. Using Wolf algorithm [25], the Lyapunov exponents spectrum of system (1) with $b_1 = 25, b_2 = 20$ and $b_4 = 7$ are represented in Figure 2.

In particular, for the parameter values as in (2), the values of Lyapunov exponents of non-linear system (1) are given by

$$L_1 = 1.968, \quad L_2 = 0.1755, \quad L_3 = -0.2914 \quad \text{and} \quad L_4 = -8.367 \quad (10)$$

3 IMPS of the identical novel hyperchaotic systems

The control goal considered in this section is that the two identical hyperchaotic systems can be achieved the IMPS. The main IMPS result via active control method is proved by using Lyapunov stability theory.

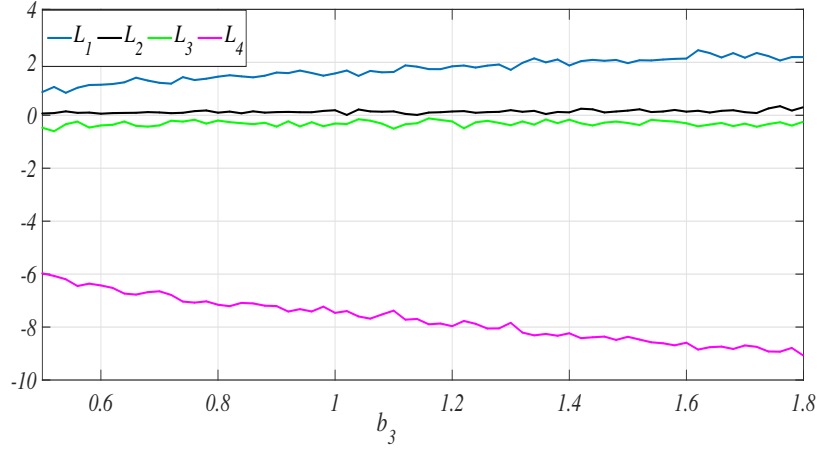


Figure 2: The Lyapunov exponents spectrum of system (1) versus b_3 .

3.1 Theoretical analysis

Consider a master and a slave hyperchaotic systems, respectively, described as follows:

$$\dot{X}(t) = f(X(t)), \quad (11)$$

$$\dot{Y}(t) = AY(t) + g(Y(t)) + U, \quad (12)$$

where $X, Y \in \mathbb{R}^n$ are state variables of the master system and the slave system, respectively, $A \in \mathbb{R}^{n \times n}$, is the linear part of the systems (12), $f, g : \mathbb{R}^n \rightarrow \mathbb{R}^n$ are non linear functions and $U = (u)_{1 \leq i \leq n}$ is a control input vector.

The problem of IMPS [15] for the systems (11) and (12) is to find the controller U such that the synchronization error,

$$e(t) = MY(t) - X(t), \quad (13)$$

satisfies

$$\lim_{t \rightarrow \infty} \|e(t)\| = 0, \quad (14)$$

where M is invertible scaling matrix. Hence, we have the following result.

Theorem 1 *The IMPS between the master system (11) and the slave system (12) can be achieved if the following conditions are satisfied*

(a) $U = M^{-1}((A - C)e(t) - MAY - Mg(Y(t)) + f(X(t)))$, where C is a control matrix.

(b) $(A - C)^T + (A - C)$ is a negative definite matrix.

Proof. The error system can be derived as follow

$$\begin{aligned} \dot{e}(t) &= M\dot{Y}(t) - \dot{X}(t) \\ &= MAY(t) + Mg(Y(t)) + MU - \dot{X}(t) \end{aligned} \quad (15)$$

Substituting (a) into Eq. (15), the error system can be described as

$$\dot{e}(t) = (A - C)e(t). \quad (16)$$

Consider the quadratic Lyapunov function defined by

$$V(t) = e^T(t)e(t), \quad (17)$$

which is positive definite on \mathbb{R}^n .

Differentiating V along the trajectories of (11) and (12), we get

$$\begin{aligned} \dot{V}(t) &= \dot{e}^T(t)e(t) + e^T(t)\dot{e}(t) \\ &= e^T(t)(A - C)^T e(t) + e^T(t)(A - C)e(t) \\ &= e^T(t) ((A - C)^T + (A - C)) e(t) \\ &< 0. \end{aligned}$$

Thus, we can conclude that all solutions of error system (13) tend towards zero exponentially as $t \rightarrow \infty$. Hence, the IMPS between the identical hyperchaotic systems (11) and (12) is achieved under the conditions (a) and (b). This completes the proof. ■

3.2 Numerical simulation

To verify the effectiveness and the feasibility of the presented synchronization method, we take the novel hyperchaotic system as a master system and its controlled system as a slave system. The master system is defined as

$$\begin{cases} \dot{x}_1 = b_1(x_2 - x_1) + x_4, \\ \dot{x}_2 = -x_1x_2 + b_2x_2, \\ \dot{x}_3 = \sinh(x_1x_2) - b_3x_3, \\ \dot{x}_4 = b_4x_1. \end{cases} \quad (18)$$

The slave system is described by

$$\begin{cases} \dot{y}_1 = b_1(y_2 - y_1) + y_4 + u_1, \\ \dot{y}_2 = -y_1y_2 + b_2y_2 + u_2, \\ \dot{y}_3 = \sinh(y_1y_2) - b_3y_3 + u_3, \\ \dot{y}_4 = b_4y_1 + u_4, \end{cases} \quad (19)$$

where u_1, u_2, u_3, u_4 are the active control functions. The linear part of the systems (18) and (19) is given by

$$A = \begin{pmatrix} -b_1 & b_1 & 0 & 1 \\ 0 & b_2 & 0 & 0 \\ 0 & 0 & -b_3 & 0 \\ b_4 & 0 & 0 & 0 \end{pmatrix}.$$

According to IMPS control technique proposed in the previous section, the matrix M and the gain matrix C are selected as

$$M = \begin{pmatrix} 2 & 0 & 0 & 2 \\ 2 & 2 & 0 & 2 \\ -1 & -1 & -1 & 0 \\ 0 & 0 & -1 & -1 \end{pmatrix}. \quad (20)$$

$$C = \begin{pmatrix} 0 & b_1 & 0 & 1 \\ 0 & 1 + b_2 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ b_4 & 0 & 0 & 3 \end{pmatrix}. \quad (21)$$

From the condition (a) of the theorem 3.1, the vector controller $u = (u_1, u_2, u_3, u_4)$ can be constructed as follows

$$\begin{cases} u_1 = \frac{1}{4}e_2 - \frac{3}{2}e_4 + x_4 - y_4 - \frac{1}{2}e_1b_1 + \frac{1}{2}e_3b_3 - b_1x_1 + b_1x_2 + b_1y_1 - b_1y_2, \\ u_2 = -\frac{1}{2}e_2 + \frac{1}{2}e_1b_1, \\ u_3 = \frac{1}{4}e_2 + \frac{3}{2}e_4 + \frac{1}{2}e_3b_3 - b_3x_3 + b_3y_3 + \sinh x_1x_2 - \sinh y_1y_2, \\ u_4 = -\frac{1}{4}e_2 + \frac{3}{2}e_4 - \frac{1}{2}e_3b_3 + b_4x_1 - b_4y_1. \end{cases} \quad (22)$$

With the choice control (22), the error system becomes

$$\begin{pmatrix} \dot{e}_1 \\ \dot{e}_2 \\ \dot{e}_3 \\ \dot{e}_4 \end{pmatrix} = \begin{pmatrix} -b_1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -b_3 & 0 \\ 0 & 0 & 0 & -3 \end{pmatrix} \begin{pmatrix} e_1 \\ e_2 \\ e_3 \\ e_4 \end{pmatrix}. \quad (23)$$

For the numerical simulations, we take:

The parameter values of the master and slave systems as in the case (2).

The initial states of the master system are taken as

$$x_1(0) = 1, x_2(0) = 1, x_3(0) = 1 \text{ and } x_4(0) = 1. \quad (24)$$

The initial states of the slave system are taken as

$$y_1(0) = 0.5, y_2(0) = 0.5, y_3(0) = 0.5 \text{ and } y_4(0) = 0.5. \quad (25)$$

We recall that the role of choice of the initial conditions is to justify the existence of hyperchaos and also for the uniqueness of the solutions of the proposed systems.

With the choice of the previous initial states, the error system has the initial states

$$e_1(0) = 1, e_2(0) = 2.5, e_3(0) = -0.5 \text{ and } e_4(0) = -2. \quad (26)$$

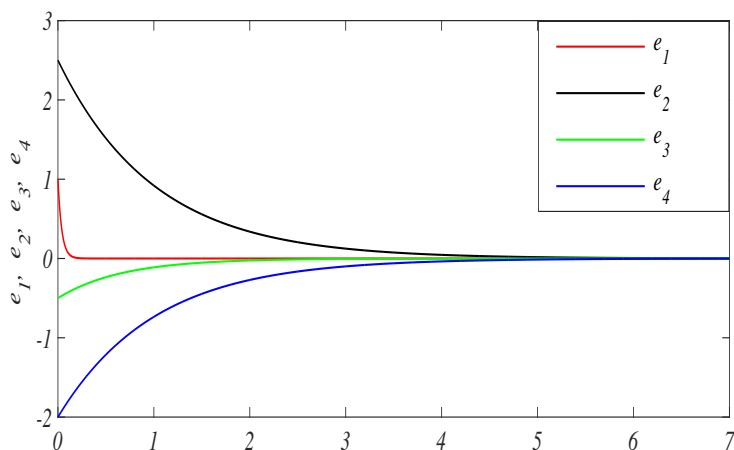


Figure 3: Time evolution of the synchronization errors (23).

The curves of synchronization error (23) are shown in Figure 3.

From Figure 3, we show that the evolution of all variables of error dynamic system (23) quickly tend towards zero as $t \rightarrow \infty$, which indicate that the IMPS between the hyperchaotic systems (18) and (19) is achieved.

4 Conclusion

In this research work, we have studied a novel hyperchaotic system with hyperbolic sine function non-linearity and their dynamic behaviours. In addition, an active controller has been proposed to ensure the possibility of IMPS of this system. Lyapunov's stability theorem has been used to provide asymptotic stability as well as convergence of synchronization errors towards zero. Finally numerical example and computer simulations are shown to illustrate all the main results of this paper.

5 Acknowledgements

This research was supported by the Algerian General Directorate for Scientific Research and Technological Development (DG-RSDT).

References

- [1] M-S. Abdelouahab, N Hamri, Fractional-order Hybrid Optical System and its Chaos Control Synchronization, *EJTP*, **30**, (2014) 49–62.
- [2] M. S. Alwan, X.Z. Liu and M.U. Akhmet, On chaotic synchronization via impulsive control and piecewise constant arguments, *Dynamics of Continuous, Discrete and Impulsive Systems, Series B: Applications and Algorithms*, **22**, (2015) 53–67.

- [3] B. Blasius, A. Huppert and L. Stone, Complex dynamics and phase synchronization in spatially extended ecological system, *Nature*, **399**, (1999) 354–359.
- [4] A. Bouzeriba, A. Boukroune and T. Bouden, Projective synchronization of two different fractional-order chaotic systems via adaptive fuzzy control, *Neural Comput. Applic.*, **27**, (2016) 1349–1360.
- [5] G.L. Cai, S. Zheng and L.X. Tian, Adaptive control and synchronization of an uncertain new hyperchaotic Lorenz system, *Chin. Phys. B*, **17**, (2008) 2412–2419.
- [6] A. Boukroune, S. Hamel, A. Bouzeriba and T. Bouden, Adaptive fuzzy control-based projective synchronization of uncertain non affine chaotic systems, *Complexity* **21**, (2014) DOI: 10.1002/cplx. 21596.
- [7] A. Chen, J. Lu, Jinhua Lü, S. Yu, Generating hyperchaotic Lü attractor via state feedback control, *Physica A*, **364**, (2006) 103–110.
- [8] T.I. Chien and T.L. Liao, Design of secure digital communication systems using chaotic modulation, cryptography and chaotic synchronization, *Chaos, Solitons and Fractals*, **24**, (2005) 241–245.
- [9] Y. Fei, H. W. Chun, H. Yan and W. Y. Jin, Anti synchronization of a novel hyperchaotic system with parameter mismatch and external disturbances, *Pramana J. Phys.*, **79**, (2012) 81–93.
- [10] M. Javidi and J.J. Nieto, Dynamic analysis of a fractional order three-level food chain model, *Dynamics of Continuous, Discrete and Impulsive Systems Series B: Applications and Algorithms*, **24**, (2017) 247–267.
- [11] M. Lakshmanan and K. Murali, *Chaos in Nonlinear Oscillators: Controlling and Synchronization*, World Scientific, Singapore 1996.
- [12] K. Murali, M. Lakshmanan, Secure communication using a compound signal from generalized chaotic systems, *Phys. Lett. A*, **241**, (1998) 303–310.
- [13] I. Ncube, Some results on the stability and bifurcation of a distribution delay network, *Dynamics of Continuous, Discrete and Impulsive Systems Series B: Applications and Algorithms*, **23**, (2016) 127–136.
- [14] A. Ouannas, S. Abdelmalek, and S. Bendoukha, Co-existence of some chaos synchronization types in fractiona-order differential equations, *Electronic Journal of Differential Equation*, **2017**, (2017) 1804–1812.
- [15] A. Ouannas, E.E. Mahmoud, Inverse matrix projective synchronization for discrete chaotic systems with different dimensions, *Journal of Computational Intelligence and Electronic Systems*, **3**, (2014) 188–192.
- [16] A. Ouannas, G. Grassi, T. Ziar, and Z. Odibat, On a function projective synchronization scheme for non-identical Fractional-order chaotic (hyperchaotic) systems with different dimensions and orders, *Optik - International Journal for Light and Electron Optics*, **136**, (2017) 513–523.
- [17] A. Ouannas, M. M. Al-sawalha, Synchronization between different dimensional chaotic systems using two scaling matrices, *Optik Int. J. for Light and Electron Optics*, **127**, (2016) 959–963.
- [18] A. Ouannas and Z. Odibat, Generalized synchronization of different dimensional chaotic dynamical systems in discrete time, *Nonlinear Dynamics*, **81**, (2015) 765–771.
- [19] A. Ouannas and Z. Odibat, On inverse generalized synchronization of continuous chaotic dynamical systems, *International Journal of Applied and Computational Mathematics*, **2**, (2016) 1–11.
- [20] A. Ouannas, A. T. Azar, V. Sundarapandian, New hybrid synchronization schemes based on coexistence of various types of synchronization between master-slave hyperchaotic systems, *Int. J. Comp. Apps. in Tech.*, **55**, (2017) 112–120.

- [21] A. Soukkou, A. Boukabou and S. Leulmi, Prediction-based feedback control and synchronization algorithm of fractional-order chaotic systems, *Nonlinear Dyn.*, **85**, (2016) 2183–2206.
- [22] A. Soukkou, A. Boukabou and S. Leulmi, Design and optimization of generalized prediction-based control scheme to stabilize and synchronize fractional-order hyperchaotic systems, *Optik*, **127**, (2016) 5070–5077.
- [23] S. Vaidyanathan, CK. Volos, VT. Pham, Analysis, adaptive control and adaptive synchronization of a nine-term novel 3-D chaotic system with four quadratic non-linearity and its circuit simulation, *J. Engi. Sci. Technol. Rev.*, **8**, (2015) 174–184.
- [24] S-G. Wang, L. Bai and M. Chen, Robust sliding mode control of general time-varying delay stochastic systems with structural uncertainties November, *Control Theory and Technology*, **12**, (2014) 357–367.
- [25] A. Wolf, J. B. Swift, H. L. Swinney, J. A. Vastano, Determining Lyapunov Exponents from a Time Series, *Physica D*, **6**, (1985) 285–317.
- [26] G. Ye and J. Zhou, A block chaotic image encryption scheme based on self-adaptive modelling, *Applied Soft Computing*, **22**, (2014) 351–357.

Received June 2018; revised January 2020.

email: smailkaouache@gmail.com